# **Annealing characteristics of Si doped amorphous silica films by rf sputtering**

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Si doped amorphous silica films were prepared by rf magnetron sputtering technique. As-deposited films which had relatively low Si dopant were annealed in inert atmosphere, and spectroscopic analyses were performed for as-deposited and all the annealed samples by infrared (IR) absorption, X-ray fluorescence analysis, X-ray photoelectron spectroscopy (XPS) and 29Si magic angle sample spinning nuclear magnetic resonance spectroscopy ( $^{29}$ Si MAS-NMR). When the sample was annealed at 800 °C, the IR absorption peak located at 1080 cm−<sup>1</sup> shifted to slightly higher wave number. On the other hand, the decrease of Si clusters or Si–Si bonds in as-deposited film was deduced from X-ray fluorescence and XPS spectra for the sample annealed at 800 °C. These annealing characteristics of the films prepared in this study were discussed based on the random bonding model of  $SiO<sub>x</sub>$  film, and the spectral variations with thermal annealing were interpreted by the rearrangement of Si and O atoms in as-deposited films, rather than the simple clustering of excess Si atoms. © 1999 Kluwer Academic Publishers

# **1. Introduction**

A silicon dioxide film on silicon has been a very useful device for a surface-protective layer, a good insulator and an antireflecting coating. Silica films prepared via vapor phase by various methods are highly pure and widely used in the electronics and optoelectronics industries. On the other hand, non-stoichiometric silicon oxide films  $(SiO_x, 0 < x < 2)$ , which are called siliconrich oxide films, are known to be interface layers between silicon substrates and silica films. SiO*<sup>x</sup>* films with various chemical compositions can be prepared by some deposition methods [1–5], and it is possible that the variation of x values causes the change in electrical and optical characteristics for  $SiO<sub>x</sub>$  films. Since  $SiO<sub>x</sub>$  films are generally annealed in the course of their processing, it is of interest to investigate the preparation process and the annealing characteristics of SiO*<sup>x</sup>* films.

Recently, Si derivative new materials, such as porous silicon and Si low-dimensional materials are of much interest for new optical devices. A new controversial phenomenon is blue photoluminescence from porous silicon after rapid thermal oxidation [6–8], which is attractive for displays and fast optoelectronics devices. The origin of this phenomenon is still unclear. Several authors have suggested that the blue photoluminescence originates from the inside of smaller Si clusters than those in red luminescent PS [6, 7]. However, other factors such as defects in a silica glass [9] or the surface states of silicon clusters [10] cannot be ruled out. Regardless of the exact origin of the photoluminescence, it may be related to the existence of Si–O chemical bonds or the oxygen-related Si microstructure. The characterization of  $SiO<sub>x</sub>$  films which have various distribution of Si–Si bonds and Si–O bonds will be more important for detailed investigations of the above Si derivative materials.

In this study, Si doped amorphous silica films were successfully prepared by the rf sputtering technique. For as-deposited and annealed samples, spectroscopic analyses were performed by infrared (IR) absorption, X-ray fluorescence analysis, X-ray photoelectron spectroscopy (XPS), and  $^{29}$ Si magic angle sample spinning nuclear magnetic resonance spectroscopy  $(^{29}Si$  MAS-NMR). The variations of these spectra with thermal annealing were discussed in terms of the change in arrangement of constituent Si and O atoms in the films.

#### TABLE I Sputtering conditions



# **2. Experimental**

Si doped amorphous silica films were prepared with rf magnetron sputtering equipment (ULVAC SH-350). The target was a pure  $SiO<sub>2</sub>$  glass plate of 4 inches in diameter on which 9 Si single-crystal chips  $(10 \times$  $10 \times 0.65$  mm<sup>3</sup>) were placed symmetrically with respect to the center. The co-sputtering was performed under the conditions of Table I. Sample films were deposited on Si single-crystal wafers for IR and XPS measurements, on polycrystalline alumina plates for X-ray fluorescence analysis and on copper substrates for 29Si MAS-NMR measurement. The film thickness was controlled by sputtering time, and it was determined from a steplike change of film edge, which was detected by a surface texture measuring instrument (Surfcom 112B Tokyo Seimitsu Corp.). For asdeposited samples, an annealing treatment was carried out under  $N_2$  3  $\times$  10<sup>2</sup> Pa atmosphere after five gas substitutions in a vacuum furnace, which was evacuated to about 3 Pa.

IR transmission spectra were recorded using a Fourier transform IR spectrophotometer (Horiba FT-310). The contribution from the silicon substrate was removed from the data using a differential method. For an X-ray fluorescence analysis for Si  $K_\beta$  line, an electron probe micro analyzer (Shimadzu EMX-SM) was used, and Si  $K_\beta$  emission spectra were excited by means of an electron beam (accelerating voltage and sample current were 20 kV and 1  $\mu$ A, respectively). XPS measurements were performed by an X-ray photoelectron spectrometer (ULVAC PHI Inc.) using a Mg  $K_{\alpha}$  X-ray source (300 W) at a constant analyzer pass energy of 35.75 eV. The calibration of the binding energy of Si 2p core level was made by fixing C 1s core level. For an NMR measurement, free standing films were prepared by stripping the films from copper substrates. A <sup>29</sup>Si MAS-NMR spectrum was measured at room temperature on a JEOL JNM-GX400 spectrometer operating at 79.3 MHz for the <sup>29</sup>Si resonance. The spinning rate was ∼6 kHz. The short pulses (1 $\mu$ s) were used with a 40 s time delay between pulses, and the spectrum was accumulated 8086 times. Chemical shifts of  $29$ Si relative to TMS (tetramethylsilane) were determined by using the peak at −34 ppm of PDMS (polydimethylsilane) as an external reference.

# **3. Results**

#### 3.1. IR absorption properties

No diffraction peak of Si microcrystals was observed for as-deposited and all the annealed samples by X-ray



*Figure 1* The variation of IR transmission spectra with annealing temperature.

diffraction measurements. As-deposited samples were 1.3  $\mu$ m and 9  $\mu$ m in thickness for 40 min and 5 h sputtering, respectively.

Fig. 1 shows the variation of IR transmission spectra with annealing temperature. The samples were annealed in  $N_2$  atmosphere for a constant annealing time, 3 h. A spectrum for silica glass film is also displayed in the figure. This silica glass film was prepared on a Si wafer by the same sputtering condition. The IR absorption peak around  $1100 \text{ cm}^{-1}$  region is related to the mode of Si–O–Si stretching vibration, and a small peak around  $800 \text{ cm}^{-1}$  is due to the Si–O–Si bending vibration [11, 12]. With the increase of annealing temperature, the main peak attributed to the Si–O–Si stretching vibration mode shifts to a slightly higher wave number. The wave number of the main peak for the silica glass film is 1080 cm−1. For as-deposited SiO*<sup>x</sup>* film, IR absorption maximum is located at  $1065 \text{ cm}^{-1}$ . From the relationship between IR peak position and x in SiO*<sup>x</sup>* [1], the x value in the  $SiO_x$  film was roughly estimated at 1.7.

# 3.2. X-ray fluorescence analysis

Fig. 2 shows the X-ray emission spectra of Si  $K_\beta$  line for a thicker sample on an alumina substrate. The thicker sample on the alumina substrate was employed for the sake of excluding the effects of the substrate. The sputtering conditions for a sample preparation are identical with that for the sample shown in Fig. 1 except for the substrate material and sputtering time. For comparison, X-ray emission spectra of Si  $K_\beta$  line for a bulk silica glass and a Si wafer are also shown with them. Si  $K_\beta$ spectra for a silica glass and a Si wafer represent electronic transitions from the valence band of silica glass to Si 1s core level and from the valence band of Si crystal to Si 1s core level, respectively [13]. The difference between Si  $K_\beta$  spectrum for as-deposited sample and that for the bulk silica glass is clearly observed around the higher energy shoulder. It would be due to



*Figure 2* X-ray emission spectra of Si  $K_\beta$  lines.

Si clusters or Si–Si bonds in the film. The shoulder of spectrum after thermal annealing is a little smaller than that of as-deposited sample. This result shows that thermal annealing makes the number of Si clusters or Si–Si bonds in as-deposited sample decrease, but a small number of Si–Si bonds are still present even after thermal annealing.

#### 3.3. X-ray photoelectron spectroscopy

Si 2p core level spectra are shown in Fig. 3. The sample shown here is the sample as shown in Fig. 1. Si 2p deeper core level energy depends greatly on the nearest neighbors of Si atoms. If some oxygen atoms of a stoichiometric  $SiO<sub>2</sub>$  glass network, where all Si atoms are surrounded by four oxygen atoms, are replaced by



*Figure 3* The variation of Si 2p core level spectra with annealing temperature.

silicon atoms that are less electronegative than oxygen atoms, the Si 2p core level will shift to lower binding energy. Rochet *et al.* [14] reported that the binding energy of the  $Si^0$  state corresponding to Si atoms having no oxygen neighbors was about 99 eV, and that of the  $Si<sup>4+</sup>$  state corresponding to Si atoms having four oxygen neighbors was about 103.5 eV.

The dashed lines in Fig. 3 represent the Gaussian curves corresponding to the  $Si<sup>4+</sup>$  state, where Si atoms are surrounded by four oxygen atoms. A slight peak of the  $Si^0$  state is observed for the as-deposited sample. Thermal annealing at 800 ◦C makes the intensity of this peak smaller, and the profile of the spectrum asymmetric, resulting from the decrease of  $Si<sup>0</sup>$  components and the increase of Si atoms in the intermediate oxidation states, such as  $Si^{2+}$  and  $Si^{3+}$ . However, annealing at  $1000\degree$ C causes the increase of the signal intensities at lower binding energies, which corresponds to the increase of  $Si^0$  and  $Si^{1+}$  components. This anomalous change of the XPS spectrum on annealing will be discussed later.

# 3.4. 29Si MAS-NMR measurement

Fig. 4 shows the  $^{29}Si$  MAS-NMR spectrum for free standing films. The sputtering conditions are identical with those for other samples with the exception of copper substrates. Sample films on copper substrates were peeled off and annealed at 800 ◦C for 3 h. The 29Si MAS-NMR chemical shift has been known to be very sensitive to the change in the first coordination sphere of the Si atom. Philipp [1] proposed that  $SiO<sub>x</sub>$ had a structure made up of Si tetrahedral units, Cn,  $Si-(O_nSi_{4-n})$ , where n took a value among 0, 1, 2, 3 and 4, and that each Si tetrahedral unit was distributed randomly in a statistical sense. From this random bonding model, Dupree *et al*. [15] estimated the peak position of the central Si atom in each tetrahedron. They performed <sup>29</sup>Si MAS-NMR measurements for three commercial SiO powders, and carried out a least-squares fit of their spectra to a sum of five Gaussians. As the result, they found that chemical shifts with the least



*Figure 4* 29Si MAS-NMR spectrum obtained at 79.3 MHz, 8086 s cans,  $1 \mu s$  pulses, 40 sec delay. Chemical shifts referenced to TMS.

uncertainty were  $-67$  and  $-109$  ppm. Referring to <sup>29</sup>Si chemical shifts of vitreous  $SiO<sub>2</sub>$  and a-Si, it was determined that these values corresponded to the chemical shifts of  $C_1$  and  $C_4$  tetrahedral units. Taking these values as fixed points and assuming a linear variation of chemical shifts, the peak position of the central Si atom in each tetrahedron was estimated. These peak positions are indicated in the figure. It is found that the sample annealed at 800 °C has  $C_2$  or  $C_3$  units in which Si atoms are in the intermediate oxidation states, but a quantitative discussion is difficult for very weak signal intensity.

#### **4. Discussion**

The change of spectra for IR absorption, X-ray fluorescence and XPS upon thermal annealing will be interpreted based on the random bonding model mentioned above. IR peak shift to higher wave number by annealing has been explained by the following two suppositions. Firstly, the macroscopic reaction

$$
2\text{SiO}_x \to x\text{SiO}_2 + (2 - x)\text{Si} \tag{1}
$$

by thermal annealing can be expected to take place. This phase separation has already been proposed [3, 14] to take place on  $SiO_x$  film with  $x < 1.5$  in the absence of oxidation. As the result of this reaction, various kinds of Si tetrahedral units in the film ultimately turn into  $C_0$  and  $C_4$  units, that is to say, Si-(Si<sub>4</sub>) and Si-(O<sub>4</sub>). Needless to say,  $C_0$  units don't contribute to Si–O–Si stretching vibration mode, thus IR peak position of this mode shifts to higher wave number. The second supposition is the oxidation of the film on thermal annealing. The oxidation of the film transforms Si tetrahedral units with small n values into the units with larger n values because of the intrusion of oxygen atoms in the film. That causes the peak shift of Si–O–Si stretching mode. In this experiment, the annealing treatment was performed under  $N_2$  3  $\times$  10<sup>2</sup> Pa atmosphere after five gas substitutions. Although the possibility of oxidation is not excluded blindly, the influence of oxidation can be almost neglected. Additionally, the anomalous change in the XPS spectrum by annealing cannot be explained by the oxidation process. Therefore, IR peak shift to higher wave numbers by annealing would be caused by the increase of  $C_4$  units through the phase separation process. However, IR spectra do not directly indicate the variation of the number of  $C_0$  units with annealing, and Si–Si vibrations cannot be detected by IR spectroscopy. The variation of the number of  $C_0$  units with annealing will be discussed below by using the other spectroscopic data.

The results of X-ray fluorescence and XPS showed the decrease of the number of  $C_0$  units by annealing at  $800\degree$ C. Without being affected by the oxidation process, this decrease would originate not from the phase separation but from the dissolution of small Si clusters. The dissolution means that  $C_0$  units in as-deposited samples turn into the units with larger n values because of their instability. The Si 2p core level spectrum for the sample annealed at 800 ℃ also indicated the decrease deposited sample. However, the Si 2p core level spectrum for the sample annealed at  $1000\degree C$  indicated the increase of the components with lower binding energy as compared with the sample annealed at  $800\degree$ C. This result supports that the change of the XPS spectrum by thermal annealing must be explained not only by the phase separation but also by the rearrangement of Si and O atoms in as-deposited film, which includes the dissolution of small Si clusters. Probably, depending on annealing conditions and on excess Si content in the film, either the phase separation or the dissolution would be dominant through the thermal annealing process. In this study, thermal annealing at temperatures below 1000 ℃ would cause the dissolution of small Si clusters, instead of the phase separation into the formation of Si microcrystals. For annealing at  $1000\,^{\circ}\text{C}$ , on the other hand, a small amount of silicon clustering would take precedence over the dissolution of Si cluster mentioned above. The increase of  $C_0$  and  $C_1$  units by annealing at  $1000\degree C$  as observed in Fig. 3 also supports a small amount of clustering of excess silicon atoms in as-deposited films.

of the number of  $C_0$  units, compared with that for the as-

Based on the random bonding theory,  $SiO_x$  film is composed of five different tetrahedral units,  $C_n$ ,  $Si-(O_nSi_{4-n})$  (*n* = 0, 1, 2, 3, 4). The rearrangement of Si and O atoms can be defined as the exchange Si atoms and O atoms between different tetrahedral units without changing the ratio of O to Si atoms. This is expressed by the equilibrium reaction,

$$
2C_n \rightleftharpoons C_{n-1} + C_{n+1} \qquad (n = 3, 2, 1). \qquad (2)
$$

On the other hand, the elementary reaction of the dissolution process would be assumed by the following reaction between five tetrahedral units,

$$
C_0 + 4C_n \to C_1 + C_{n-1} + 3C_n \qquad (n = 3, 2). \tag{3}
$$

If the equilibrium reaction of Equation 2 ( $n = 3$ ) proceeds to the right and the reaction of Equation 3 takes place, the increase of  $C_4$  units and the decrease of  $C_0$ units by annealing are achieved simultaneously. If only a small number of  $C_0$  units exist in the as-deposited film, this rearrangement will be the origin of the IR peak shift to higher wave number for the sample annealed at 800 ℃. The rearrangement into the dissolution of small Si clusters (Equation 3) would cause the depression of the shoulder for the X-ray fluorescence spectrum by annealing at 800 °C and the decrease of  $Si^0$ and  $Si<sup>1+</sup>$  components by annealing at the temperature below  $1000\,^{\circ}\text{C}$  as shown in Figs 2 and 3. This model for the rearrangement of Si and O atoms gives a satisfactory explanation for the spectral variation of the film with fewer excess Si atoms by annealing.

Surely, this supposition based on a simple model will not be perfectly consistent with the real structural change of the sample. However, this is a possible mechanism of an annealing process for Si doped amorphous silica films with fewer excess Si atoms. The result of the <sup>29</sup>Si MAS-NMR measurement, where the existence of  $C_2$  units rather than  $C_0$  units were observed, also

shows that thermal annealing at 800 ◦C for this sample does not cause the macroscopic phase separation. The phase separation and the dissolution process by thermal annealing successfully account for annealing characteristics of Si doped amorphous silica films. For more quantitative understanding of the annealing characteristics of  $SiO_x$  films with x value approaching 2, theoretical calculations and thermodynamic studies will be required.

#### **5. Conclusion**

Si doped amorphous silica films were prepared by the rf magnetron sputtering technique, using 9 Si chips on a silica glass target. For as-deposited and annealed samples, spectroscopic analyses were performed by IR absorption, X-ray fluorescence, XPS and 29Si MAS-NMR. When the sample was annealed up to 1000 ◦C, IR absorption peak shifted to a slightly higher wave number. On the other hand, the binding energy of the Si 2p core level varied anomalously with annealing temperature. These spectral variations of the films on thermal annealing were interpreted by the rearrangement of Si and O atoms in as-deposited films. For Si doped amorphous silica films with fewer excess Si atoms, thermal annealing at temperatures below 1000 ◦C would cause the dissolution of small Si clusters, rather than the phase separation into the formation of Si microcrystals.

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#### **References**

- 1. H. R. PHILIPP , *J. Phys. & Chem. Solids* **32** (1971) 1935.
- 2. E. HOLZENKÄMPFER, F.-W. RICHTER, J. STUKE and U. VOGET-GROTE, *J. Non-Cryst. Solids* **32** (1979) 327.
- 3. L. A. NESBIT, *Appl. Phys. Lett*. **46** (1985) 38.
- 4. S. HAYASHI, T. NAGAREDA, Y. KANZAWA and K. YAMAMOTO, *Jpn. J. Appl. Phys*. **32** (1993) 3840.
- 5. Y. OSAKA, K. TSUNETOMO, F. TOYOMURA, H. MYOREN and K. KOHNO, *ibid*. **31** (1992) L365.
- 6. M. K. LEE and K. R. PENG, *Appl. Phys. Lett*. **62** (1993) 3159.
- 7. H. MIMURA, T. FUTAGI, T. MATSUMOTO, T. NAKAMURA and Y. KANEMITSU, *Jpn. J. Appl. Phys.* **33** (1994) 586.
- 8. L. TSYBESKOV, JU. V. VANDYSHEV and P. M. FAUCHET, *Phys. Rev. B* **49** (1994) 7821.
- 9. A. G. CULLIS, L. T. CANHAM, G. M. WILLIAMS, P. W. SMITH and O. D. DOSSER, *J. Appl. Phys.* **75** (1994) 493.
- 10. K. TAKEDA and K. SHIRAISHI, *Solid State Commun.* **85** (1993) 301.
- 11. M. NAKAMURA, Y. MOCHIZUKI, K. USAMI, Y. ITOH and T. NOZAKI, *Solid State Commun.* **50** (1984) 1079.
- 12. P. G. PAI, S. S. CHAO, Y. TAKAGI and G. LUCOVSKY, *J. Vac. Sci. & Technol. A* **4** (1986) 689.
- 13. D. L. GRISCOM, *J. Non-Cryst. Solids* **24** (1977) 155.
- 14. F. ROCHET, G. DUFOUR, H. ROULET, B. PELLOIE, J. PERRIERE, E. FOGARASSY, A. SLAOUI and M. FROMENT, *Phys. Rev. B* **37** (1988) 6468.
- 15. R. DUPREE, D. HOLLAND and D. <sup>S</sup> . WILLIAMS , *Philos. Mag. B* **50** (1984) L13.

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